OXIDE OR CARBIDE NANOPARTICLES SYNTHESIZED BY LASER ABLATION OF A BULK Hf TARGET IN LIQUIDS AND THEIR STRUCTURAL, OPTICAL AND DIELECTRIC PROPERTIES

N. G. Semaltianos^{*a}, J. -M. Friedt^b, R. Chassagnon^e, V. Moutarlier^d, V. Blondeau-Patissier^b, G. Combe^b, M. Assoul^c, G. Monteil^c

^aAristotle University of Thessaloniki, Dept. of Physics, Thessaloniki, 54124, GREECE

^bUniversité de Franche-Comté, Femto-st, Dépt. Temps-Fréquence, UMR CNRS 6174, Besançon, 25030, FRANCE

^cUniversité de Franche-Comté, Femto-st, Dépt. Mécanique Appliquée, UMR CNRS 6174, Besancon, 25030, FRANCE

^dUniversité de Franche-Comté, UTINAM, UMR CNRS 6213, Besançon, 25030,

FRANCE

^eUniversité de Bourgogne Franche-Comté, Laboratoire Interdisciplinaire Carnot De Bourgogne, ICB UMR CNRS 6303, Dijon, 21078, FRANCE

Abstract

Laser ablation of a bulk Hf target in deionized (DI) water, ethanol or toluene was carried out for the production of nanoparticles colloidal solutions. Due to the interaction of the ablation plasma plume species with the species which are produced by the liquid decomposition at the plumeliquid interface, hafnia (HfO₂) nanoparticles are synthesized in DI water, hafnium carbide (HfC) nanoparticles in toluene and a mixture of those in ethanol. The hafnia nanoparticles are in the monoclinic low temperature phase and in the tetragonal and fcc high temperature phases. Their size distribution follows log-normal function with a median diameter in the range of 4.3-5.3 nm. Nanoparticles synthesized in DI water have band gaps of 5.6 and 5.4 eV, in ethanol 5.72 and 5.65 eV (using low and high pulse energy) and in toluene 3 eV. Values for the relative permittivity in the range of 7.74-8.90 were measured for hafnia nanoparticles thin films deposited on substrates by drop-casting (self-assembled layers), in parallel plate capacitor structures.

Keywords: laser ablation; nanoparticles; hafnia; hafnium carbide

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*Corresponding author, E-mail: nsemaltianos@yahoo.com

1. Introduction.

Hf is a tetravalent transition metal element, in the same column in the periodic table as Zr and Ti. In particular, hafnium dioxide (HfO₂) termed as hafnia and other hafnia-based materials find wide applications in the field of microelectronics as good replacements of silica in field effect transistor gates due to their good balanced combination of properties such as a high dielectric constant, relatively large band gap, large heat of formation, good thermal and chemical stability on Si, large barrier heights at interfaces with Si and adequate compatibility with n^+ polysilicon gate electrodes [1]. Hafnia is used as a high refractive index material and in the form of thin films as heat resistant and strongly reflective protective optical coating [2]. Hafnia nanoparticles find applications as luminescent nanomaterials [3], in the medical field for nanoscale radiotherapy of tumor cells by achieving localized intracellular high energy gamma/x-ray radiation dose deposit (due to their high-Z number) [4], in the fabrication of nanocomposites with polymers to be used as γ -ray scintillators [5], as fillers in proton conducting membranes for proton-exchange membrane fuel cells, in high etch resistance lithographic photoresists for EUV patterning [6] and others.

The most common traditional methods of synthesis of hafnia nanoparticles in solution, involve the reaction of hafnium(IV) ethoxide with benzyl alcohol [7], the reaction of hafnium tetrachloride (HfCl₄) with potassium hydroxide (KOH) and 3,4-dihydroxy hydrocinnamic acid (DHCA) [8] or just with sodium hydroxide (NaOH) [9], the reaction of hafnium(IV) isopropoxide with HfCl₄ and trioctylphosphine oxide (TOPO) as the strongly coordinating solvent [10] or the reaction of hafnium(IV) tert-butoxide with ammonia (NH₄OH) in the presence of surfactant [11]. These methods require long reaction times and high temperatures [7,8,10], annealing at usually high temperatures under controlled gas atmospheres of the as synthesized originally amorphous nanoparticles in order to crystallize them [9,11] or vigorous

washing procedures in order to remove from the nanoparticles' surfaces the coordinating ligands (such as for instance TOPO) which are used during their synthesis [10].

In addition to hafnia, another interesting material from the hafnium compounds is the hafnium carbide (HfC). This is the most refractory binary compound known with a melting point over 3890 °C, high phase stability and good thermomechanical and thermochemical properties. It has attracted investigations in recent years for ultra-high temperature applications such as in cutting tools, rocket nozzles, space/air craft and thermal-field emitters [12]. It is usually synthesized by the reaction of powder mixtures of some source of carbon such as phenolic resin, citric acid monohydrate, graphite powder and others, with some source of Hf such as pure Hf metal, Hf hydride (HfH₂), hafnia or HfCl₄ [11-15].

An alternative method for the synthesis of nanoparticles in solution involves laser ablation of a bulk target material while the material is immersed in a liquid [16-18]. The nanoparticles are formed by the debris which is generated as a result of the material's ablation and a nanoparticles colloidal solution is formed. In case of ablation in water, ultrapure nanoparticle colloidal solutions are produced, without any reaction by-products since the process of nanoparticles synthesis does not involve the use of any chemical precursors and most importantly nanoparticles with bare, ligand-free surfaces are synthesized. This is the main reason why this method has gained popularity and especially in the fabrication of polymernanoparticles nanocomposites since the interaction of the nanoparticles with their environment mainly depends on the properties of the nanoparticles' surface. The produced colloidal solutions are stable against nanoparticles' agglomeration without the need of adding into the solution any stabilizing surfactants, due to the electrostatic repulsion between the nanoparticles because of the surface electrical charge which is generated onto them from their surface partial oxidation and subsequent ionic equilibrium with the liquid [19,20]. In addition, the ablation plasma plume experiences an additional compression by the liquid which may result in the formation of nanoparticles which are characterized by metastable material phases, difficult or impossible to be produced by other methods. Throughout the years, nanoparticles out of a number of different materials have been synthesized using this method [16-18].

Laser ablation (120 fs, 800 nm, 1 kHz, 0.1-0.36 mJ) of a Zr target in ammonia or water lead to the formation of zirconia (ZrO_2) nanoparticles in the cubic or tetragonal/monoclinic phase, respectively [21]. H⁺ and Al²⁺- codoped Al₂O₃ nanoparticles (rather than aluminium hydrates) with spinel related structures of mainly γ - and its monoclinic derivative θ -phase having significant internal compressive stress and defects which exhibit band gap (5.2 eV) much lower than the band gap of α -Al₂O₃ (8.8 eV) were synthesized by laser ablation (16 ns, 1064 nm, 10 Hz, 850 mJ) of an Al target in DI water [22] while ablation in aqueous solution of NaOH (0.05-1 M) lead to the fabrication of epitaxial β -NaAlO₂ and γ -Al₂O₃ nanoparticles [23]. Hollow amorphous Al₂O₃ micro/nanoparticles which may had weak crystallization were synthesized by pulsed KrF excimer laser ablation (30 ns, 248 nm, 10 Hz, 2.3 J/cm²) of a bulk Al target in distilled water [24]. Laser ablation (90 fs, 800 nm, 1 kHz, 110 µJ) of a Cr target in DI water lead to the formation of a mixture of Cr₃O₄, Cr₂O₃ and CrO₃ nanoparticles, in ethanol to a mixture of Cr₃O₄, Cr₇C₃ and of the metastable phase Cr₃C_{2-x} nanoparticles while in acetone or toluene to nanoparticles which were in the metastable phase Cr₃C_{2-x} [25]. W₂C-C core-shell nanoparticles or pearl-like SiC nanorings were synthesized by laser ablation (10 ns, 1064 nm, 20 Hz, 100 mJ) of a W or Si target, respectively in ethanol [26]. Laser ablation (Nd:YAG laser, 1064 nm, 10 Hz, 90 mJ) of a Ta target in ethanol lead to the synthesis of Ta_xO crystal core/Ta₂O₅ amorphous shell nanoparticles [27].

This paper involves the synthesis of nanoparticles in solution by laser ablation for the first time of a bulk Hf target in DI water, ethanol or toluene. The morphological, structural, optical and dielectric properties of the nanoparticles were measured and compared with the corresponding properties of a similar type of nanoparticles synthesized in the literature by other

methods. The synthesized nanoparticles exhibit different crystal phases and the underlying physics for the formation of the different phases in relation to the temperature and pressure conditions of the plasma plume during its adiabatic expansion in the liquid environment, is discussed. In addition to the band gap absorption, sub-band optical transitions are also observed in the UV-vis absorption spectra and their presence is explained based on the possible formation of defects in the nanoparticles or of any deviations from the perfect stoichiometry. Finally, a general methodology is outlined for the measurement of the dielectric properties of nanoparticles-based thin films.

2. Experimental Details.

Laser ablation for nanoparticles generation was carried out in this work using a Q-switched Nd:YAG pulsed laser (pulse width=20 ns) at beam wavelength of 1064 nm and pulse repetition rate of 1 kHz using pulse energy of 380 or 1194 μ J. The target was a bulk piece of Hf metal (purity 99.5 %) with thickness of 0.75 mm and dimensions of 10×10 mm. Before each ablation the target surface was mechanically polished using SiC paper in order to remove any surface oxide layers and expose a clean material's surface. The target was used for ablation immediately after polishing and ablation was carried out with the target lying at the bottom of a beaker filled with 5 ml of the liquid. The height of the liquid above the target surface was ~10 mm. The beam was impinging vertically onto the horizontally lying target at the bottom of the beaker, scanning the target surface in a meander fashion (5mm length, pitch 0.02 mm, scanning speed of 1 mm/s, 400 overscans) and ablation was carried out for ~15 mins each time resulting to ~1-1.3 mg of nanoparticles in the 5 ml of each liquid. Formation of the nanoparticles resulted in the initial solvents becoming from transparent to white in the case of ablation in DI water and brownish in the case of ablation in ethanol or toluene (Fig. S1 [64]). The nanoparticles colloidal solutions were characterized by UV-vis spectrophotometry. The synthesized nanoparticles, from solution

droplets dried out onto carbon coated copper grids, were characterized by Transmission Electron Microscopy (TEM) and from droplets dried out onto clean glass substrates by X-Ray Diffraction (XRD) using a diffractometer with a Cu K_{α} source (λ =1.5406 Å). Samples for Fourier Transform Infrared (FTIR) transmission measurements were prepared using the KBr method, by mixing dried nanoparticles with KBr (in a ratio of ~8:168 mg) and preparing a pellet with thickness of ~1 mm. Crystal structures were determined from the XRD patterns using the software of the instrument (DiffracPlus EVA).

Samples for capacitance measurements were prepared by depositing the nanoparticles as a thin film onto Au/Cr (100/20 nm thickness) metallized by RF-sputtering, circular glass substrates (thickness ~0.5 mm, diameter 10 mm) by drop casting from the colloidal solutions while keeping the substrates in an oven at temperature of ~60 °C. After completion of the deposition, the samples were kept in the oven at the same temperature for a further ~ 48 hrs, for a final drying. In order to avoid any contamination of the samples, the whole process was performed in a clean room environment. This method provides with a self-assembled layer of nanoparticles which has a quite uniform thickness over the whole substrate surface area (variation of ~10 %) of 13.72 ± 1.15 µm (data for other samples are reported in Table 1 of the Supplementary Material [64]). The thickness was measured by physically scratching the film surface and measure the depth of the scratches by a white light scanning interferometric microscope. Then a Au film (thickness ~100 nm) was deposited onto the nanoparticles layer, by RF-sputtering, using a metallic mask with circular holes with diameters of 0.5, 1.0, 1.5, 2.0 and 2.5 mm (Fig. S3 [64]). The metallic mask was fabricated from an 80 µm thick stainless steel foil, by laser cutting. In this way circular plate capacitors with different diameters were fabricated. The admittance versus frequency of the capacitors was measured in the range 0.1-10 MHz using a vector network analyzer from which the dielectric constant of the nanoparticles film was determined.

3. Results and Discussion.

3.1 Synthesis of nanoparticles and characterization of their morphological, structural and optical properties.

The laser beam with fluence greater than the ablation threshold of the material produces at the place of incidence on the material's surface a plasma plume of ablated species (neutral atoms or ions). This plasma plume in the case of long pulse laser ablation as at the present case, where the deposition of energy in the material is slow, originates from vaporization of the material melt pool, namely the ablation of the material is via melting and vaporization rather than by sublimation (direct solid to vapour transition) as in the case of ultrashort pulse laser ablation [25]. The nanoparticles are formed by nucleation in the vapour phase (of the plume species) but a large percentage of nanoparticles correspond to material which is directly ejected from the material's surface in the form of molten droplets [28].

TEM images of the nanoparticles are shown in Fig. 1 (a)-(d). It is seen that spherical nanoparticles are generated in DI water and ethanol while in toluene there are also a few nanoparticles which are not perfectly spherical. This is most probably caused by the low temperature conditions of the plasma plume relative to the high temperature which is required for the reaction of Hf with C in comparison to the one which is required for its reaction with oxygen species. These conditions result to the reduced mobility of surface atoms during nuclei coalescence and thus to the imperfect smoothing of grain boundaries. These nanoparticles are formed by nuclei whose temperature is not high enough to allow the transport of interface atoms during approximately 400 nanoparticles in images of nanoparticles ensembles obtained on different areas on each grid, are shown in Fig. 1 (e)-(h). The histograms are described quite well by lognormal functions with median diameters of $<d_0>=4.6$, 5.3, 4.5 and 4.3 nm and standard deviations of $\sigma=0.55$, 0.56, 0.52 and 0.45, respectively.

XRD patterns measured from the nanoparticles are shown in Fig. 2. The as measured patterns are presented without any smoothing of the curves applied for noise reduction. In the same graphs the reference patterns for the several material structures which are identified in the nanoparticles are also plotted for comparison. Analysis of the patterns using the software for the diffractometer reveals that the peaks observed in the patterns for the nanoparticles synthesized in DI water (Fig. 2 (a)) correspond to hafnia in the monoclinic, tetragonal and fcc phase [29]. The monoclinic is an equilibrium phase of hafnia and is observed at normal conditions of pressure and temperature. The tetragonal and fcc are high temperature phases of hafnia and are observed at temperature above ~1700 °C and ~2500 °C, respectively, also at low pressure [30]. In traditional method of synthesis of hafnia nanoparticles these phases can be stabilized at much lower temperatures but by the addition into hafnia of rare earth dopants such as Y, Yb, Gd and others [31,32] or Si [33]. The origin of the peaks at 36.646°, 48.203° and 55.687° which are observed in the patterns for the nanoparticles synthesized in DI water (labeled with the symbol #) is still unknown, but they might be related to the orthorhombic phase of hafnia with space group Pbca termed as OrthoI which was determined from neutron diffraction studies [34]. This is a high pressure phase of hafnia which is observed in the pressure range of 4-14 GPa [30]. Formation of hafnia nanoparticles is due to the interaction (oxidation) of the hafnium plasma plume species with the oxygen species which are produced at the plume-liquid interface by the thermolysis of water because of the extreme temperature of the plume. The presence in the synthesized nanoparticles of the crystal phases with different stabilization conditions of temperature and pressure is explained by the highly non-equilibrium conditions of the plasma plume. At the start of the expansion of the plume when its temperature is high, the fcc phase is formed and then as the plume expands and adiabatically cools, the tetragonal and monoclinic phases are formed. The high pressure which is required for the formation of the orthorhombic phase is provided by the "compression" which the plume experiences by the liquid during its adiabatic expansion in the liquid confined environment [16-18]. When later the plume expansion quenches the high temperature and pressure phases are "frozen" in the generated nanoparticles. Interestingly, althought the peaks in the pattern for the nanoparticles synthesized in DI water using higher pulse energy of 1194 μ J become stronger in intensity as compared to the peaks for the nanoparticles synthesized using lower pulse energy of 380 μ J (due to the higher quantity of material produced by the ablation), their relative intensities ratio remains approximately the same (Fig. S2 [64]). This indicates that the relative percentage of each crystal phase (monoclinic, tetragonal and fcc) in each nanoparticles ensemble is approximately the same, independent of the pulse energy used at least in the energies range used in this work.

The peaks in the pattern for the nanoparticles synthesized in ethanol (Fig. 2 (b)) correspond to hafnia in the monoclinic, tetragonal and fcc phases and also to HfC [29,35]. The peaks in the pattern for the nanoparticles synthesized in toluene (Fig. 2 (c)) correspond to HfC while the peak at 45.494° to pure Hf with hexagonal lattice [36]. The oxide nanoparticles in the case of ablation in DI water or ethanol are formed by the interaction of the plasma plume species with the oxygen species which are produced by the liquid decomposition at the plume-liquid interface [25]. The carbide nanoparticles are formed by the interaction of the plume species with the carbon from the pyrolysis of the liquid (carburization of the plume species). The presence of pure Hf nanoparticles in addition to the HfC ones is due to the incomplete carburization of the plume species due to the insufficient amount of carbon atoms relative to the amount of Hf species available for reaction as well as to the relatively high temperature which is required for the reaction of the metal with carbon as compared to the one which is required for its reaction with oxygen, as mentioned previously. These are mainly the larger nanoparticles which are ejected directly from the target material as melted droplets.

UV-vis absorption spectra from the nanoparticles colloidal solutions are shown in Fig. 3. Considering a direct band gap, parabolic valence and conduction bands and negligible reflectance at low wavelengths, the band gap of the nanoparticles can be found from the plot of $(\alpha\omega)^2$ versus energy ($\hbar\omega$) by the interception of the extrapolation of the linear part of the curve (at high energies) with the horizontal (energy) axis ($\hbar\omega=0$) (Fig. 4). Small deviation from the linearity might be caused by the fact that tetragonal hafnia has an indirect band gap [37]. Band gap values of 5.6 and 5.4 eV, are obtained for the nanoparticles synthesized in DI water using low and high pulse energy, respectively (Fig. 4(a)). A large variation of the band gap values of hafnia nanoparticles or thin films exists in the literature. A value of 5.55 eV was measured for pure hafnia [38]. A value of 6.02 eV was reported for the band gap of hafnia nanoparticles in the monoclinic phase with an average diameter of $\sim 15.5\pm 2$ nm [9]. A band gap value of 5.65 eV was reported for 500 nm thick hafnia thin films in the monoclinic phase [39]. Band gap values in the range of 5.85-5.90 eV were reported for monoclinic hafnia thin films with thicknesses between 5 and 40 nm while lower values in the range of 5.61-5.80 eV were reported for mixed monoclinic, orthorhombic and tetragonal phase films [40]. In mixed monoclinic and cubic phase yttria-doped hafnia thin films the band gap was observed to increase with the percentage of the cubic phase as the film thickness was increasing, from 5.60 (± 0.03) eV for pure monoclinic phase films to 5.70-5.85 $(\pm 0.04 \text{ eV})$ for mixed films to 6.05 (±0.03) eV for pure cubic phase films [41]. Band gaps in the region of 5.50-5.85 eV were obtained for hafnia thin films which were thermally annealed in the 500-1000 °C temperature range, assuming indirect transitions [42]. The higher value for the band gap which is obtained for the nanoparticles synthesized using a lower energy as compared to the value for the band gap for the nanoparticles synthesized using a higher energy might be due to the quantum confinement effect because of the smaller median diameter of the nanoparticles (4.6 as compared to 5.3 nm, Fig. 1(e) and (f)), similarly to the case of hafnia thin films (2, 3 or 4 nm thickness) to which the increase of the band gap was attributed to the quantum confinement effect along the vertical direction [43] in analogy with quantum well heterostructures.

For the nanoparticles synthesized in ethanol a band gap value of 5.65 eV is determined (and 5.72 eV for the nanoparticles synthesized using lower energy) (Fig. 4 (b)). However it is seen that there are also weak bands in the absorption spectra (Fig. 3 (b)) peaked at around 215 and 240 nm which result in an energy value of 4.95 and 3.8 eV, respectively (5.1 and 3.85 eV for the nanoparticles synthesized using low energy) (Fig. 4 (b)). Such bands with energy lower than the band gap of the material have also been observed previously in hafnia thin films [44,40] and were found to originate from defects states produced by the distortion of bonds at the internal boundaries of the nanocrystalline grains which are present in the nanoparticles (grain-boundary defect states) [45] or from possible oxygen vacancies due to a deviation of the elemental stoichiometry of the nanoparticles from that of hafnia's [40]. These states act as bulk traps in advanced Si devices and are deleterious for the device reliability and performance [45]. The distortion of bonds might be expressed as an increase or decrease of bond lengths in different crystal directions. This is observed at the present work in the XRD patterns (Fig. 2) in which the angles at which few of the peaks corresponding to a particular phase of the material appear, are not exactly equal to the values of the angles at which those peaks are expected to appear according to the corresponding reference tables (a deviation of up to $\sim 0.5\%$). The peaks appear at slightly higher or lower angles indicating a shrinking or an expansion of the unit cell, respectively along the corresponding crystal direction.

For the nanoparticles synthesized in toluene a band gap energy of 3 eV is obtained (Fig. 4 (c)) which is close to the band gap energy measured for HfC powders using EELS spectroscopy [46]. A small deviation might be caused from the presence of also Hf in the nanoparticles ensemble (in addition to HfC).

FTIR transmission spectra from hafnia nanoparticles synthesized in DI water are shown in Fig. 5. The bands in the low wavenumbers region at 610, 748 and 642 cm⁻¹ are phonon modes of monoclinic and tetragonal hafnia, respectively [47-49]. The band at 1587 cm⁻¹ is due

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to Hf-OH [47] while the band at 3410 cm⁻¹ is a O-H stretching band which corresponds to OH groups either bonded to H in molecular water or to Hf (Hf-OH) [50]. This indicates the presence of OH groups on the nanoparticles surfaces and confirms the absence of any other ligands on their surfaces. The band at 1630 cm⁻¹ is a H-OH bending vibration of molecular water [47] but it may also correspond to Hf-H bonds [51] while the bands at 1022, 1354, 1420, 2854 and 2928 cm⁻¹ originate from C-O vibrations in CO₂ or C-H vibrations in hydrocarbons from the ambient atmosphere [50] as well as all the remaining unmarked peaks in the spectrum.

3.2 Dielectric properties of hafnia nanoparticles.

A schematic diagram of the structure of the samples which were fabricated for the investigation of the dielectric properties of hafnia nanoparticles is shown in Fig. 6 [52]. The equivalent circuit is a parallel plates capacitor where the top and bottom plates are the gold thin films and the dielectric between the capacitor plates consists of the nanoparticles self assembled layer.

The measured Re*Y* and Im*Y* of the admittance versus frequency (Supplementary Material [64]) are shown in Fig. 7 (a) and (b) for the four devices with different surface areas. From the slopes of the curves of Im*Y* versus *f* (Fig. 7 (b)) the capacitance (*C*') versus area (*A*) was determined and is plotted in Fig. 7 (d). From the slope of the line and taking into account the film thickness, the relative permittivity of the nanoparticles film was determined as: $\varepsilon_r' = 7.74\pm1.40$. The slopes of the curves of Re*Y* versus *f* determined from Fig. 7 (a) versus the area (*A*) are plotted in Fig. 7 (c) and from the slope of the line the imaginary part of the relative permittivity was determined as: $\varepsilon_r' = 0.32\pm0.07$ Thus the loss tangent was determined as: $\tan \delta = 0.042\pm0.001$. Values for the relative permittivity and loss tangent for other samples are reported in Table 1 and Fig. S5 and S6 (Supplementary Material) [64].

The value for the relative permittivity which is obtained here is by almost three times lower than the value reported for bulk hafnia ($\varepsilon_{HfO_2} = 25$) [1]. A range of different values for the relative permittivity of hafnia have been reported in the literature including also relatively low values. Based on the numerous reports in the literature on the subject, it is concluded that in general the relative permittivity of hafnia thin films depends on the crystal structure of the film but also on the film deposition method (wet chemical or dry physical or chemical methods) [1]. The relative permittivity of uniform and continuous hafnia thin films deposited by atomic layer deposition (ALD) was measured equal to ~ 20 [53]. A value of 16 was reported for mixed monoclinic and tetragonal phase thin films [54]. In studies of the relative permittivity of hafnia nanopowder (pure moniclinic phase) the permittivity of relatively thick layers (1.5-2.5 mm thickness) was measured to take the value of ~10.5 at 10^{-2} Hz and 180 °C [55]. Hafnia thin films deposited by Molecular Beam Epitaxy were measured to have dielectric constant of around 19-20 [56]. A value of 7.5 was measured for amorphous hafnia thin films deposited on Si substrates by pulsed laser deposition (PLD) while upon annealing the films at 500 °C, either in air or nitrogen atmosphere, the structure was transformed to monoclinic or orthorhombic and with higher values of the dielectric constant of 40.8 and 25.1, respectively [57]. The increase of the dielectric constant was attributed to both the densification of the film morphology caused by the heating and to the change of the crystal structure. Higher values for the dielectric constant in the region of 17-24 were measured for as deposited hafnia thin films, deposited by PLD on Si substrates heated at 700 °C [58]. Calculations have shown that the relative permittivity of hafnia very much depends on its crystal structure with values reported as 26.17 for the cubic phase, 28.54 for the tetragonal phase and 15.47 for the monoclinic phase [59]. Low values of the relative permittivity in the range of 5.6-6.4 were also measured for hafnia thin films of thickness 10 or 30 µm almost independent on the film thickness, which were deposited on substrates by PLD (laser intensity=10¹⁰ W/cm²) [60]. The crystal structure

of hafnia was a mixture of monoclinic, tetragonal and cubic [61]. By mechanically compressing the films the permittivity was measured to take higher values in the range of 11.7-12.7 and thus the low value of the permittivity for the as deposited films was attributed to the presence of air voids in the gaps between the nanoparticles (the as deposited films had 60 ± 5 % porosity which was decreasing to 20 ± 5 % for the compressed films). If we assume that the low value for the permittivity which we obtain here is due to the presence of air-voids in the films (Fig. 6) then using the Effective Medium Approximation (EMA) [62,60] we can estimate the percentage of porosity for the films studied here. According to the EMA the dielectric constant of mixed phase materials is described by the following relation:

$$f_a \frac{\varepsilon_a - \varepsilon}{\varepsilon_a + 2\varepsilon} + f_b \frac{\varepsilon_b - \varepsilon}{\varepsilon_b + 2\varepsilon} = 0$$
 [Eq. 1]

with

$$f_a + f_b = 1$$
 [Eq. 2]

where ε_a , ε_b and f_a , f_b are the dielectric constants of the two phases a and b and their relative percentage composition, respectively, and ε is the dielectric constant of the effective medium (nanoparticles+air). By taking into account that: $\varepsilon \equiv \varepsilon'_r = 7.74$ and $\varepsilon_b \equiv \varepsilon_{air} = 1$, then if we consider $\varepsilon_a \equiv \varepsilon_{HIO_2} = 25$ equal to the value for bulk hafnia [1], the percentage of voids is estimated equal to 51 %. Data for the other samples are given in Table 1 (Supplementary Material) [64]. Although mechanical compression of the nanoparticles thin films leading to an increase of the measured dielectric constant was adopted as a method for confirming the presence of air voids in the film as responsible for the measured low value of the dielectric constant, however compression of the film might have an obvious effect in the shape and size/size distribution of the nanoparticles or even in their crystal structure. For instance, previous studies have shown that the interface strain of hafnia nanoparticles thin film with the underlying substrate causes the presence of cubic and tetragonal phases in addition to the monoclinic phase [63,1]. Compression of the nanoparticles film might induce a mechanical stress in the nanoparticles and also lead to a deformation of the shape of the nanoparticles and/or causing the nanoparticles to effectively join together in the form of a bulk material thus resulting to a change of the crystal structure of the nanoparticles as that was originally formed by the synthesis conditions. Hafnia thin films are typically used in device fabrication and other applications as they are deposited, since compression or any other attempt for a post-deposition modification of the morphology or structure of the film might interfere with the device structure and cause unpredictable errors, and according to the literature the dielectric constant of the film systematically depends not only on the crystal structure of the material but also on the method of deposition. In this work the nanoparticles were deposited on the substrate by drop-casting from the colloidal solution, forming a self-assembled layer and approximately the same value for the dielectric constant within experimental error was measured for a number of samples which were deposited under the same conditions. This indicates a consistency of the measurements in relation to the set experimental conditions and establishes the value for the dielectric constant which is measured here as the value which is obtained by adopting this particular method for the deposition of the nanoparticles thin films, that of drop-casting.

4. Conclusions.

In this work the laser ablation products of a bulk Hf target in liquids were investigated. The interaction of the plasma plume species with the oxygen or carbon species originating from the thermal decomposition of water, ethanol or toluene lead to the direct, one-step synthesis of HfO_2 or HfC nanoparticles. Due to the laser ablation synthesis the nanoparticles have bare, ligand-free surfaces which is advantageous for their use in high k-dielectric structures, in the densification process of ceramics by hot pressing or plasma sintering, in the

fabrication of polymer-nanoparticle nanocomposites or other applications. The high temperature fcc and tetragonal hafnia phases in addition to the monoclinic phase are also present in the ablation products which is highly desirable in applications as these phases are known to have higher dielectric constants than the monoclinic phase while using traditional methods, complicated, resources and time consuming processes are usually required for their synthesis.

Supplementary Material

See supplementary material for additional information.

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[64] See supplemental material at [URL] for the parallel plates capacitor model, bottles with the nanoparticles colloidal solutions, rescaled XRD data for the nanoparticles synthesized in DI water using the different energies, gold top electrodes and capacitance data for additional samples including a data summarizing Table.

FIGURE CAPTIONS

FIG. 1. TEM images and corresponding size distribution histograms of the nanoparticles synthesized in: DI water (a), (b) and (e), (f), ethanol (c) and (g) (1194 μ J), and toluene (d) and (h) (380 μ J), respectively.

FIG. 2. XRD patterns of the nanoparticles.

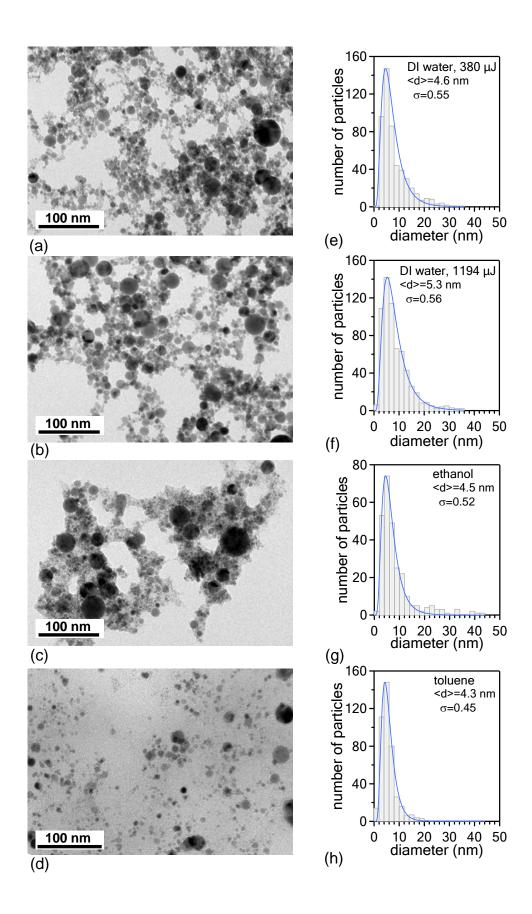
FIG. 3. UV-vis absorption spectra measured from the nanoparticles colloidal solutions (as measured).

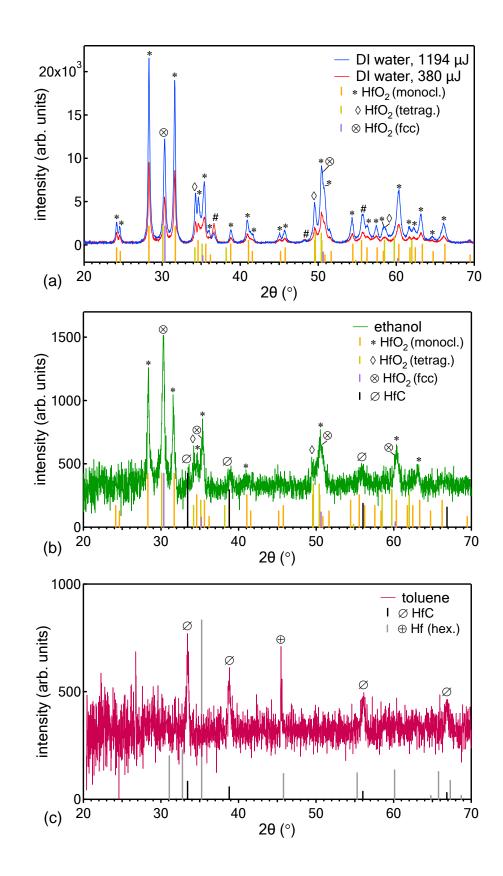
FIG. 4. Plots of $(\alpha\omega)^2$ versus energy calculated from the spectra of Fig. 3, from which the energy gaps were determined.

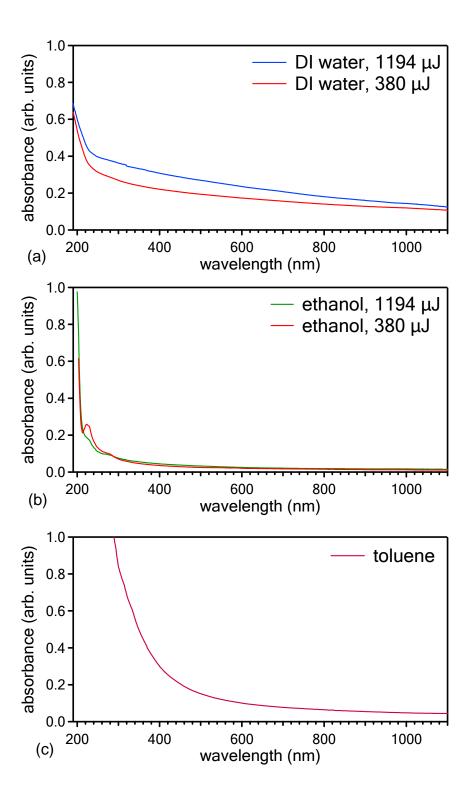
FIG. 5. FTIR transmission spectrum from the hafnia nanoparticles (1194 μ J in DI water).

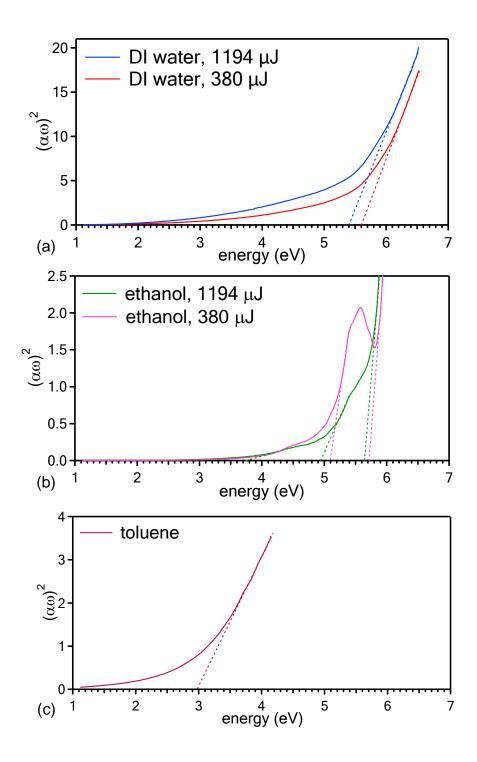
FIG. 6. Schematic of the structure fabricated for the measurement of the dielectric properties of hafnia nanoparticles thin film deposited on the substrate by drop-casting from the colloidal solution (self-assembled layer).

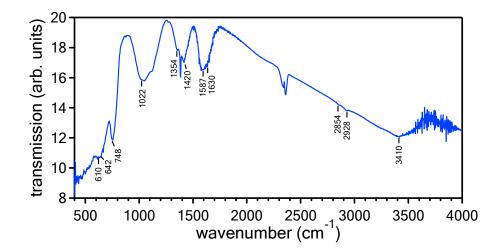
FIG. 7. Real (a) and imaginary (b) part of the admittance versus frequency, and plots of the slope of the real part (c) and of the capacitance (d) versus the device area for sample 1. Data for other samples are given in the Supplementary Material (Table 1 and Fig. S5 and S6 [64]).

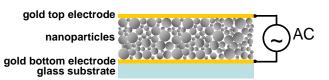


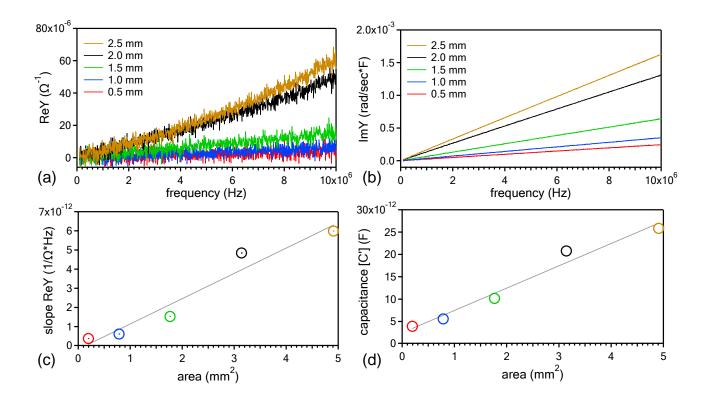












OXIDE OR CARBIDE NANOPARTICLES SYNTHESIZED BY LASER ABLATION OF A BULK HF TARGET IN LIQUIDS AND THEIR STRUCTURAL, OPTICAL AND DIELECTRIC PROPERTIES

N. G. Semaltianos^{*a}, J. -M. Friedt^b, R. Chassagnon^e, V. Moutarlier^d, V. Blondeau-Patissier^b, G. Combe^b, M. Assoul^c, G. Monteil^c

^aAristotle University of Thessaloniki, Dept. of Physics, Thessaloniki, 54124, GREECE

^bUniversité de Franche-Comté, Femto-st, Dépt. Temps-Fréquence, UMR CNRS 6174,

Besançon, 25030, FRANCE

^cUniversité de Franche-Comté, Femto-st, Dépt. Mécanique Appliquée, UMR CNRS 6174, Besançon, 25030, FRANCE

^dUniversité de Franche-Comté, UTINAM, UMR CNRS 6213, Besançon, 25030,

FRANCE

^eUniversité de Bourgogne Franche-Comté, Laboratoire Interdisciplinaire Carnot De

Bourgogne, ICB UMR CNRS 6303, Dijon, 21078, FRANCE

*Corresponding author, E-mail: nsemaltianos@yahoo.com

Parallel plates capacitor model.

The capacitance of the parallel plates capacitor with lossy dielectric (Fig. S4 (a)) is given by the relation [1]:

$$C = \frac{\varepsilon_0 \varepsilon_r A}{d}$$
 [Eq. S1]

where $\varepsilon_0 = 8.854 \times 10^{-12} F/m$ is the vacuum permittivity, $A = \pi r^2$ is the area of the capacitor plate, *d* is the distance between the plates (equal to the film thickness) and ε_r is the relative permittivity of the material between the capacitor plates which in general is a complex function:

$$\varepsilon_r = \varepsilon_r' - j\varepsilon_r''$$
 [Eq. S2]

where the real part ε_r' represents the relative permittivity (static dielectric contribution) in capacitance calculation and the imaginary part ε_r'' represents the energy loss in the dielectric medium. Furthermore, the loss tangent which is defined as:

$$\tan \delta \equiv \frac{\varepsilon_{r}}{\varepsilon_{r}}$$
 [Eq. S3]

represents how lossy the material is for AC signals.

The admittance is given by the relation:

$$Y \equiv \frac{1}{Z} = j\omega C = \frac{\omega\varepsilon_0 \varepsilon_r^{'} A}{d} + j \frac{\omega\varepsilon_0 \varepsilon_r^{'} A}{d}$$
[Eq. S4]

where Z is the impedance of the circuit. The admittance of the dielectric medium can be considered equivalent to the parallel combination of an ideal lossless capacitor C' with a relative permittivity ε'_r and a conductance G = 1/R, where R is the resistance (Fig. S4 (b)). Therefore the admittance can be written as:

$$Y = G + j\omega C'$$
 [Eq. S5]

Thus from Eqs. S4 and S5 the real ($\operatorname{Re} Y$) and imaginary ($\operatorname{Im} Y$) part of the admittance are:

$$\operatorname{Re} Y \equiv G = \frac{\omega \varepsilon_0 \varepsilon_r^{"} A}{d} = \frac{2\pi \varepsilon_0 \varepsilon_r^{"} A}{d} f \quad \text{and} \quad \operatorname{Im} Y \equiv \omega C' = 2\pi f C' = \frac{2\pi \varepsilon_0 \varepsilon_r^{"} A}{d} f \quad [\operatorname{Eq. S6}]$$

References

[1] J. M. Herbert, "Ceramic Dielectrics and Capacitors", CRC Press, (1985).

FIGURE CAPTIONS

FIG. S1. Nanoparticles colloidal solutions by laser ablation of the bulk target material in: (a1) toluene, and (a2, inset) DI water.

FIG. S2. XRD patterns from the nanoparticles.

FIG. S3. Gold top electrodes on the nanoparticles films of the fabricated capacitor structures.

FIG. S4. Equivalent circuit of the structure of Fig. 6, with a capacitor: (a) with lossy dielectric, and (b) with lossless dielectric.

FIG. S5. Real (a) and imaginary (b) part of the admittance versus frequency, and plots of the slope of the real part (c) and of the capacitance (d) versus the device area for sample 2.

FIG. S6. Real (a) and imaginary (b) part of the admittance versus frequency, and plots of the slope of the real part (c) and of the capacitance (d) versus the device area for sample 3.

Table 1. Dielectric measurements on sample 1 and on additional samples 2 and 3. The symbol (\neq) indicates a short circuit.

sample	1		2		3	
έr	7.74±1.40		8.90±1.73		8.84±1.13	
$\mathcal{E}_{r}^{"}$	0.32±0.07		0.48±0.12		0.29±0.05	
tanδ	0.042±0.001		0.054±0.003		0.032±0.002	
thickness (µm)	13.72±1.15		12.47±1.20		13.60±0.83	
% air voids	51		47		48	
diameter	Ċ	slope ReY	Ċ	slope ReY	Ċ	slope ReY
(mm)	(F)	$(\Omega^{-1} \text{Hz}^{-1})$	(F)	$(\Omega^{-1} \text{Hz}^{-1})$	(F)	$(\Omega^{-1} \text{Hz}^{-1})$
0.5	3.9068e-12	3.7922e-13	4.2822e-12	4.1171e-13	4.5711e-12	4.8276e-13
1.0	5.5792e-12	6.0919e-13	6.1014e-12	7.1115e-13	7.1036e-12	1.0931e-12
1.5	1.0173e-12	1.5259e-12	¥	¥	1.4366e-11	2.7784e-12
2.0	2.0771e-11	4.8420e-12	2.2256e-11	6.4081e-12	2.1026e-11	3.8123e-12
2.5	2.5832e-11	5.9898e-12	¥	¥	¥	¥

